## **ASTR 421** Stellar Observations and Theory

## Lecture 05 Opacity

1

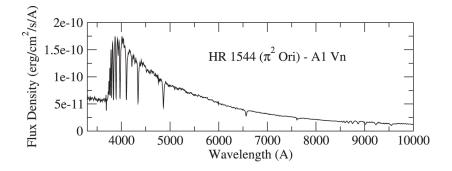
Prof. James Davenport (UW)

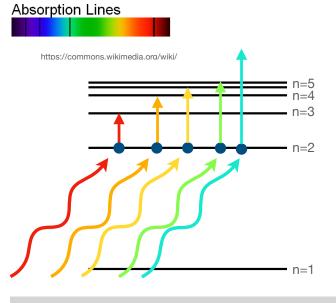
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- Boltzmann & Saha Eqns
- Properties of real spectra (e.g. continuum, lines, jumps/breaks)





**Continuous Spectrum** 

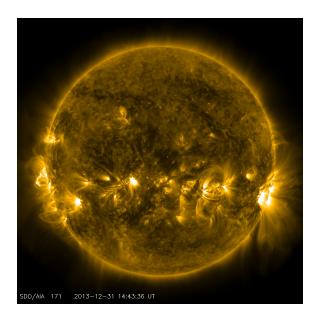
**Emission Lines** 

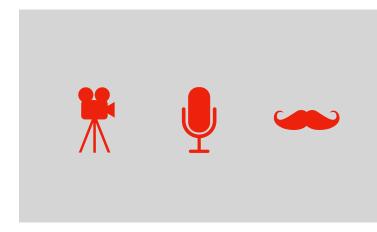


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#### **Today's Goal: Understanding Opacity**

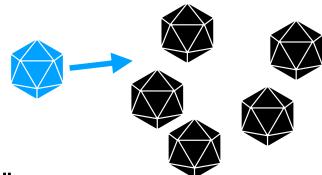
- How is opacity defined?
- How do we define the "surface" of a star
- Types of opacity in a star





#### **Mean Free Path**

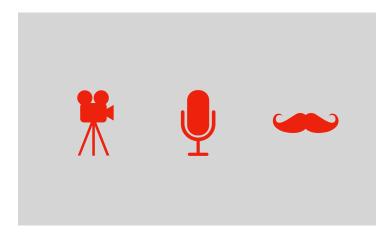
- How far can a particle travel before colliding with another?  $l = -\frac{1}{l}$ 
  - $n\sigma$
- *n* is the "number density" [1/cm<sup>3</sup>]



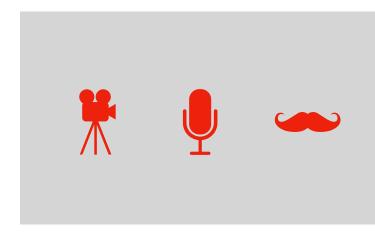
4

- $\sigma$  is the "collisional cross section" [cm<sup>2</sup>]
- Important concept for gas dynamics (e.g. Maxwell-Boltzmann velocity distribution) and for "opacity" (i.e. photons interacting with things)

- Definition: ability for material to **absorb or scatter** a photon
  - Relatedly: how far can a photon travel in a material before it is absorbed/scattered



- Definition: ability for material to **absorb or scatter** a photon
  - Relatedly: how far can a photon travel in a material before it is absorbed/scattered
- High opacity (i.e. opaque), photons can't make it through the object
- Low opacity (i.e. transparent), photons make it through w/ some absorption.
- Primarily due to: density (actually the "column density")
- Also due to: **composition, temperature**



- Opacity, or absorption coefficient, usually written as:  $\kappa$
- $[\kappa] = \mathrm{cm}^2/\mathrm{g}$

looks like: (cross-section per particle) / (g of material)

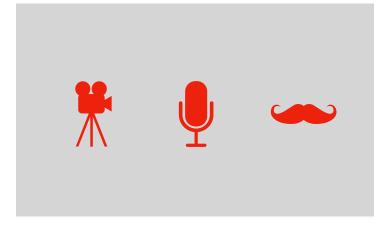
Metals have high opacity!

• This is why it's "column density" that determines observed opacity.





7



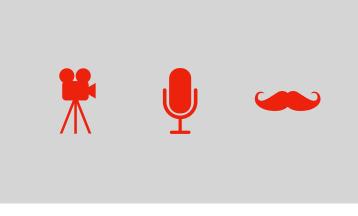
cm<sup>2</sup>

• Change of intensity over a distance (s) due to an opaque material is:  $dI_{\lambda} = -\kappa_{\lambda}\rho I_{\lambda}ds$ 



- So,  $I_{\lambda} = I_{(\lambda,0)} e^{-\kappa_{\lambda}\rho s}$ 
  - $\kappa\rho$  is like the "length scale" for absorption here
  - Recall mean free path for collisions

$$l = \frac{1}{n\sigma} = \frac{1}{\kappa\rho}$$



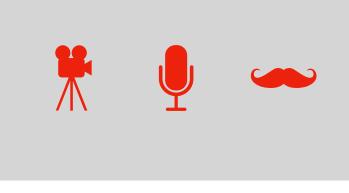
#### **Optical Depth**

- Not the same as the characteristic length!
- Important concept for stars, defining where emission comes from

9

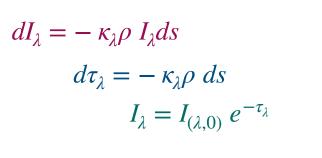
 $dI_{\lambda} = -\kappa_{\lambda}\rho \ I_{\lambda}ds$  $d\tau_{\lambda} = -\kappa_{\lambda}\rho \ ds$  $I_{\lambda} = I_{(\lambda,0)} \ e^{-\tau_{\lambda}}$ 

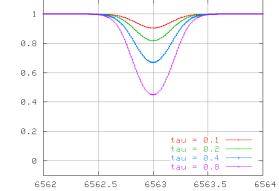
- Where  $\tau$  (the optical depth) is the number of mean free paths along a line of sight, or the "thickness" of a gas
- $\tau >> 1$ : optically thick,  $\tau << 1$ : optically thin



#### **Optical Depth**

• In simplest case, weak lines give you direct map of optical depth for the optically thin absorbing material



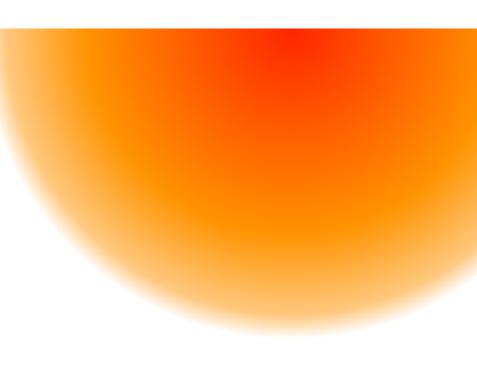


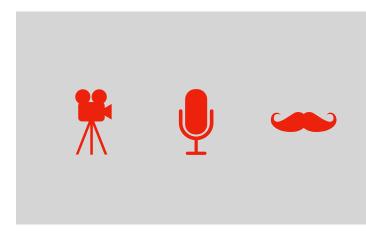
http://spiff.rit.edu/classes/phys440/lectures/curve/curve.html



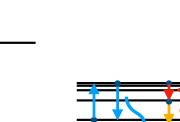
#### Surface of a Star

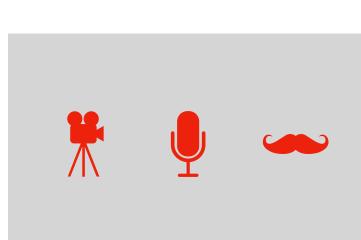
- $\tau \equiv 0$  at the surface
- We see light down to  $\tau = 1$ 
  - This is heavily wavelength dependent! (even in spectral lines themselves!)
- Photosphere is classically defined as:  $\tau_{\lambda} \approx 2/3$ 
  - Photosphere is wavelength dependent!

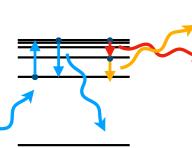




- <u>Absorption</u>: photon is removed (for a while)
  - <u>Bound-Bound</u>, at discrete wavelengths, forms absorption lines, as discussed. If quickly re-emits, acts like scattering, & can change the  $\lambda$ Sometimes called "line scattering"
  - <u>Bound-Free</u>, above a wavelength threshold, photoionization

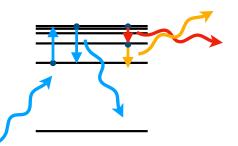


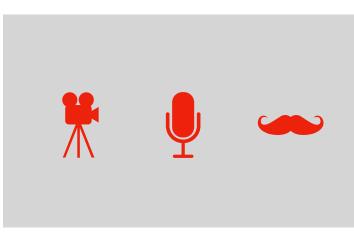


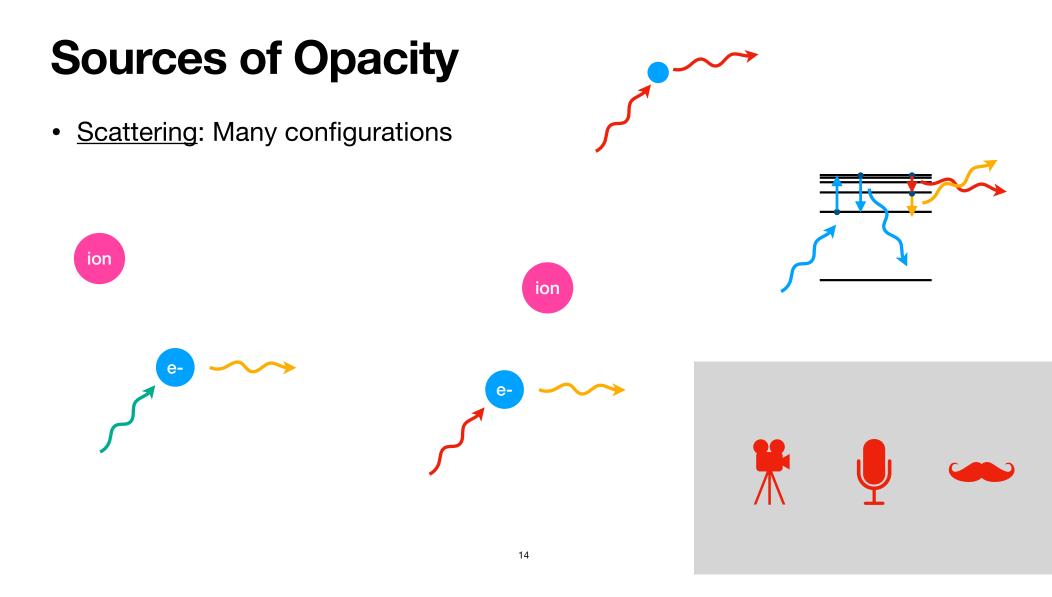




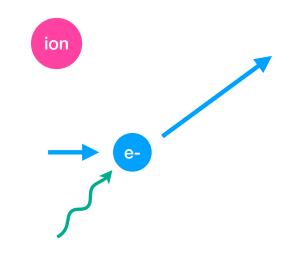
- Scattering: photon is redirected
- Acts as absorption in one direction & emission in another
- Many kinds of scattering (names can be confusing) work in a similar way:
  - Photon approaches electron
  - Electron "vibrates" due to photon's B field
  - Photon leaves in another direction





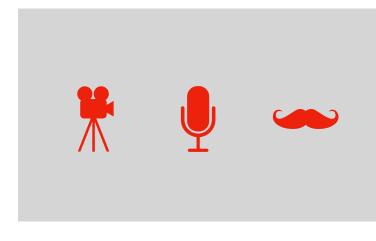


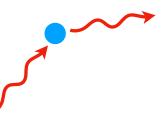
- Scattering:
- Free-Free: absorption by a "free" electron in the presence of (but not bound to) an ion
- Can occur from any wavelength photon, acts as a "continuous" opacity
- The opposite scenario: an e- loses energy passing by an ion, emits a photon, called "Bremsstrahlung" radiation.
  - This includes both "synchrotron" & "cyclotron" radiation, but we won't discuss here



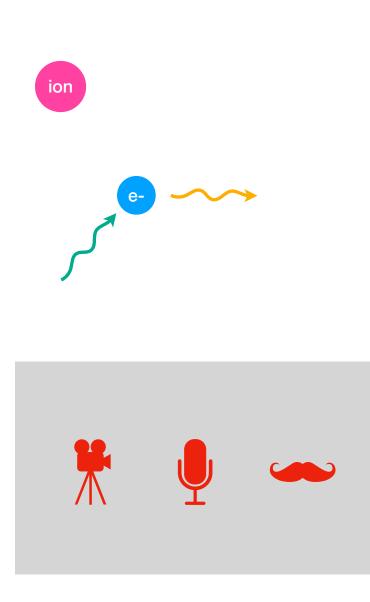


- Scattering:
- Pure (elastic) electron scattering is called Thompson Scattering
- Cross-section is related to the classical radius of the electron!  $\sigma_T = \frac{8\pi}{3} r_e^2 \approx 6.65\text{e-}29 \text{ m}^2$
- Not wavelength dependent, so only depends on *electron density*
- Important for very hot (ionized, lots of e-) and dense gas: deep in stellar interiors!



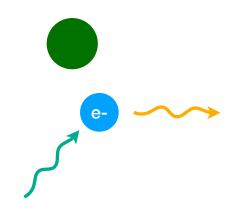


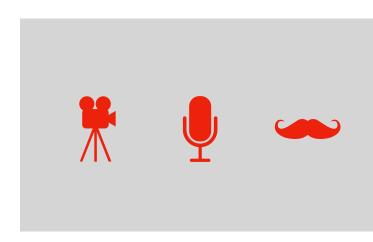
- Scattering:
- Compton Scattering: inelastic
- If the e- is in the presence of an ion, the e- can exchange momentum from the photon
- This usually results in heating of the e-(lower energy photon)
- Can also impart energy to the photon, called the *inverse* Compton effect
- Relativistic effects are important for detailed calc.



- Scattering:
- Similar to Compton, **Rayleigh Scattering** happens often for electrons around molecules or atoms
- Very wavelength dependent  $\sigma \sim 1/\lambda^4$
- This is broadly the cause of the blue sky, & the colors at sunset

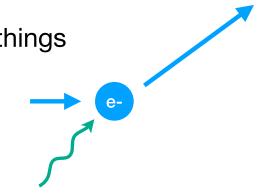






#### **Radiation Pressure (Briefly)**

- Photons have momentum, and so exert pressure on things (called radiation pressure)  $p = h/\lambda$
- If you had TOO much light, this pressure would start launching material off a star
- This is known as the **Eddington Luminosity**
- More opacity, easier to impart pressure from radiation





#### **Next time:**

- Opacity of the H- ion in the Sun
- Limb Darkening
- Total Opacity
- Spectral Lines
- Reading: BOB Ch 9.1-9.2; LeBlanc Ch 3.1-3.3

